

# Condensed Matter - HW 10 :: Superfluid

PHSX 545

## Problem 1

Using the phonon-roton model of  $^4\text{He}$  spectrum, find the quasiparticle entropy, heat capacity and density of the normal component at low temperature.

## Problem 2

(a) Find the energy  $E$  and the angular momentum  $L$  carried by a single vortex placed along the rotation axis of a long cylinder of radius  $R$ . Take the superfluid wave function in the form  $\psi(r, \phi) = \sqrt{N/V} e^{in\phi} \theta(r - r_c)$  with the uniform superfluid density  $\rho_s$  ( $n = 0, 1, 2, \dots$  and  $\theta(x)$  is the step-function). From this deduce the critical angular frequency  $\omega_c$  for nucleation of a single vortex (when  $E_{rot} = E - L\omega$  becomes negative).

(b) A neutron star of radius  $R = 10\text{km}$  rotates with frequency  $\Omega/2\pi = 1\text{ Hz}$ . Find the area density  $n_v$  of vortices in neutron superfluid, assuming the rotational motion of the star is due to vortices, their density is uniform, and the single-quantized vortices are along the rotation axis. Hint: calculate the circulation along a circle of radius  $r$  inside the star.

### Answer of exercise 1

The energies and distribution functions of the phonons and rotons in the reference frame of (non-moving) condensate are obtained from the Hamiltonian  $\mathcal{H} = \sum_{p \neq 0} \varepsilon_p \hat{b}_p^\dagger \hat{b}_p$  with:

$$\begin{aligned} \varepsilon_{ph}(p) &= up & n_{ph}(p) &= \frac{1}{e^{\beta u p} - 1} \\ \varepsilon_r(p) &= \Delta + \frac{(p-p_0)^2}{2m_r} & n_r(p) &\approx e^{-\beta \varepsilon_r(p)} \end{aligned}$$

The partition function and free energy of the excitations is

$$Z = \text{Tr} \{ e^{-\beta \mathcal{H}} \} = \prod_{p \neq 0} \frac{1}{1 - e^{-\beta \varepsilon_p}} \quad \Rightarrow \quad F = -T \ln Z = T \sum_{p \neq 0} \ln(1 - e^{-\beta \varepsilon_p}) = - \sum_{p \neq 0} \varepsilon_p - T \sum_{p \neq 0} \ln n_p$$

The entropy carried by the quasiparticles can be brought to the standard form using the Bose distribution function,

$$S = - \frac{\partial F}{\partial T} = \sum_{p \neq 0} [(n_p + 1) \ln(n_p + 1) - n_p \ln n_p]$$

and from it we get the heat capacity:

$$C_v = T \left( \frac{\partial S}{\partial T} \right)_v = \sum_{p \neq 0} \varepsilon_p \frac{\partial n_p}{\partial T} = \frac{\partial E}{\partial T}$$

We can calculate entropy directly, or we can start by calculating the energy. The phonon contribution:

$$\text{Energy:} \quad E_{ph} = \int \frac{d^3 k}{(2\pi\hbar)^3} \frac{uk}{e^{\beta uk} - 1} = \frac{u}{2\pi^2 \hbar^3} \int_0^\infty \frac{k^3 dk}{e^{\beta uk} - 1} = \frac{T^4}{2\pi^2 \hbar^3 u^3} \int_0^\infty \frac{x^3 dx}{e^x - 1} = \frac{T^4}{2\pi^2 \hbar^3 u^3} \underbrace{\Gamma(4)}_{3!} \underbrace{\zeta(4)}_{\pi^4/90}$$

$$\text{Specific heat:} \quad \boxed{C_{ph} = \frac{\partial E}{\partial T} = T^3 \frac{4\Gamma(4)\zeta(4)}{2\pi^2 \hbar^3 u^3}} \quad - \text{ now agrees with experimental } T^3!$$

$$\text{Entropy:} \quad \boxed{S_{ph} = \int dT \frac{C}{T} = \frac{1}{3} T^3 \frac{4\Gamma(4)\zeta(4)}{2\pi^2 \hbar^3 u^3} = \frac{C_{ph}}{3}}$$

The roton contribution:  $\Delta \gg T$  and  $n_p \ll 1$ :

$$\begin{aligned} \text{Energy:} \quad E_r &= \int \frac{d^3 k}{(2\pi\hbar)^3} \left( \Delta + \frac{(k-p_0)^2}{2m_r} \right) e^{-\frac{\Delta}{T} - \frac{(k-p_0)^2}{2m_r T}} = \frac{e^{-\frac{\Delta}{T}}}{2\pi^2 \hbar^3} \int_0^\infty k^2 dk \left( \Delta + \frac{(k-p_0)^2}{2m_r} \right) e^{-\frac{(k-p_0)^2}{2m_r T}} \\ &= \frac{e^{-\frac{\Delta}{T}}}{2\pi^2 \hbar^3} \int_{-\infty}^\infty (k+p_0)^2 dk \left( \Delta + \frac{k^2}{2m_r} \right) e^{-\frac{k^2}{2m_r T}} \approx \frac{p_0^2 \Delta \sqrt{2\pi m_r T}}{2\pi^2 \hbar^3} e^{-\frac{\Delta}{T}} \end{aligned}$$

$$\text{Specific heat:} \quad \boxed{C_r = \frac{\partial E}{\partial T} = \frac{p_0^2 \Delta \sqrt{2\pi m_r T}}{2\pi^2 \hbar^3} \left( \frac{1}{2T} + \frac{\Delta}{T^2} \right) e^{-\frac{\Delta}{T}}}$$

$$\text{Entropy directly from definition:} \quad S_r \approx \sum_{p \neq 0} [n_p - n_p \ln n_p] = \sum_{p \neq 0} n_p \left( 1 + \frac{\varepsilon_r(p)}{T} \right) \approx \frac{\Delta}{T} \frac{e^{-\Delta/T} p_0^2}{2\pi^2 \hbar^3} \int_0^\infty dp e^{-\frac{(p-p_0)^2}{2m_r T}}$$

$$\boxed{= \frac{p_0^2 \Delta \sqrt{2m_r \pi T}}{2\pi^2 \hbar^3 T} e^{-\Delta/T}}$$

The density of normal component (quasiparticles) for Bose particles is

$$\rho_n(T) = \frac{1}{3} \int \frac{d^3 p}{(2\pi\hbar)^3} p^2 \left( -\frac{dn_p}{d\varepsilon_p} \right) = \frac{1}{3} \frac{1}{2\pi^2 \hbar^3 T} \int_0^\infty p^4 dp \frac{e^{\beta \varepsilon_p}}{(e^{\beta \varepsilon_p} - 1)^2} = \frac{1}{3} \frac{1}{2\pi^2 \hbar^3} \begin{cases} \frac{T^4}{u^5} 4\Gamma(4)\zeta(4) & , \quad \text{phonons} \\ \frac{p_0^4 \sqrt{2\pi m_r T}}{T} e^{-\Delta/T} & , \quad \text{rotons} \end{cases}$$

### Answer of exercise 2

(a) The energy of a vortex line with phase of the order parameter  $\theta(\mathbf{r}) = n\phi$  in the center of the cylindrical vessel of height  $H$  is

$$E = H \int_0^{2\pi} d\phi \int_{r_c}^R r dr \frac{\rho_s v_s^2}{2} \quad \mathbf{v}_s = \frac{\hbar}{m} \nabla \theta = \frac{\hbar n}{m} \hat{\phi}$$

that gives for energy of a single vortex

$$E_1 = H 2\pi \frac{\rho_s n^2 \hbar^2}{2m^2} \ln \frac{R}{r_c}$$

The angular momentum of a vortex is

$$L_1 = H \int d\phi \int_0^R r dr (r \rho_s v_s) = 2\pi H \rho_s \frac{\hbar n R^2}{2m}$$

and the energy in the rotating reference frame

$$E_{rot} = E - L\omega$$

becomes negative when a single vortex nucleation is favorable:

$$\omega > \omega_c = \frac{E_1}{L_1} = \frac{\hbar n}{m R^2} \ln \frac{R}{r_c}$$

(b) We can write the velocity field due to many vortices as a sum over individual vortices at positions  $\mathbf{r}_i$ :

$$\nabla \times \mathbf{v}_s(\mathbf{r}) = \sum_i 2\pi \frac{\hbar}{m} \delta(\mathbf{r} - \mathbf{r}_i)$$

The mass here is the mass of particles making up Bose condensate. In neutron star this is pairs of neutrons (2 fermions making up a Bose particle), so  $m = 2m_n$ . The circulation along circle of radius  $r$  is integration over area of the contour:

$$\oint_C \mathbf{v}_s \cdot d\mathbf{r} = \iint (\nabla \times \mathbf{v}_s) \cdot d\mathbf{S} = 2\pi \frac{\hbar}{m} N_C$$

where  $N_C$  is the number of vortices inside the circle,  $N_C = n_v \pi r^2$ . Assuming azimuthal symmetry the contour integral on the left is  $2\pi r v_s(r)$  and the velocity of the superfluid flow as a function of the radius is

$$\mathbf{v}_s(r) = \frac{\hbar}{m} \frac{n_v \pi r}{2} \hat{\phi} \quad - \text{ just like solid body rotation } \omega r$$

and thus we have

$$\Omega = \frac{\hbar \pi n_v}{2m} = \frac{\hbar n_v}{4m} \quad \Rightarrow \quad n_v = \frac{4m\Omega}{\hbar} = \frac{4 * 2 * 1.7 \cdot 10^{-27} \text{ kg} * 1 \text{ Hz}}{10^{-34} \text{ m}^2 \text{ kg Hz}} \sim 10^8 \text{ 1/m}^2$$

or 10,000 per square centimeter. Vortex separation distance is about 0.1 mm.